

FlaviA
net



Chiral low-energy constants from τ data

Martín González-Alonso
martin.gonzalez@ific.uv.es

Instituto de Física Corpuscular
(CSIC – Universidad de Valencia)



Contents

- L_{10} & C_{87} : Motivation.
- Approach: sum rules with the V-A correlator.
- Data side.
- ChPT side.
- Results & comparisons.
- Conclusions.

(M. G.-A., A. Pich & J. Prades, work in progress)

L_{10} & C_{87} : Motivation.

- ChPT Lagrangian:

$$\mathcal{L}^{\chi PT} = \frac{F^2}{4} \langle u_\mu u^\mu + \chi_+ \rangle + \dots + \frac{1}{4} L_{10}^r \langle f_{+\mu\nu} f_+^{\mu\nu} - f_{-\mu\nu} f_-^{\mu\nu} \rangle + \dots +$$

$$+ \dots + C_{87}^r \langle \nabla_\rho f_{-\mu\nu} \nabla^\rho f_-^{\mu\nu} \rangle + \dots + O(p^8)$$

- The motivation is two-fold:

- **Phenomenology:**

more precise LEC's \rightarrow more precise predictions.

For example...

$$\pi \rightarrow e \nu \gamma$$

- **Theoretical:**

Our estimation (directly from the data) allows us to test the quality of the different theoretical models that have predicted these LEC's.

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

- V-A correlator:

$$\begin{aligned}\Pi_{V-A}^{\mu\nu}(q) &\equiv i \int d^4x e^{iqx} \langle 0 | T(J_L^\mu(x) J_R^\nu(0)^\dagger) | 0 \rangle \\ &= (q^\mu q^\nu - q^2 g^{\mu\nu}) \Pi^{(0+1)}(q^2) + q^2 g^{\mu\nu} \Pi^{(0)}(q^2)\end{aligned}$$

$$\begin{aligned}J_L^\mu &= \bar{u} \gamma^\mu (1 - \gamma_5) d \\ J_R^\mu &= \bar{u} \gamma^\mu (1 + \gamma_5) d\end{aligned}$$

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

- V-A correlator:

$$\begin{aligned}\Pi_{V-A}^{\mu\nu}(q) &\equiv i \int d^4x e^{iqx} \langle 0 | T(J_L^\mu(x) J_R^\nu(0)^\dagger) | 0 \rangle \\ &= (q^\mu q^\nu - q^2 g^{\mu\nu}) \Pi^{(0+1)}(q^2) + q^2 g^{\mu\nu} \Pi^{(0)}(q^2)\end{aligned}$$

$$\begin{aligned}J_L^\mu &= \bar{u} \gamma^\mu (1 - \gamma_5) d \\ J_R^\mu &= \bar{u} \gamma^\mu (1 + \gamma_5) d\end{aligned}$$

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

- V-A correlator:

$$\begin{aligned}\Pi_{V-A}^{\mu\nu}(q) &\equiv i \int d^4x e^{iqx} \langle 0 | T(J_L^\mu(x) J_R^\nu(0)^\dagger) | 0 \rangle \\ &= (q^\mu q^\nu - q^2 g^{\mu\nu}) \Pi^{(0+1)}(q^2) + q^2 g^{\mu\nu} \Pi^{(0)}(q^2)\end{aligned}$$

$$\begin{aligned}J_L^\mu &= \bar{u} \gamma^\mu (1 - \gamma_5) d \\ J_R^\mu &= \bar{u} \gamma^\mu (1 + \gamma_5) d\end{aligned}$$

- Weight function: $w(z)\Pi(z)$

$$w(s) = \frac{1}{s}, \frac{1}{s^2}$$

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

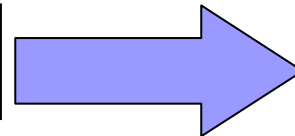
- V-A correlator:

$$\begin{aligned}\Pi_{V-A}^{\mu\nu}(q) &\equiv i \int d^4x e^{iqx} \langle 0 | T(J_L^\mu(x) J_R^\nu(0)^\dagger) | 0 \rangle \\ &= (q^\mu q^\nu - q^2 g^{\mu\nu}) \Pi^{(0+1)}(q^2) + q^2 g^{\mu\nu} \Pi^{(0)}(q^2)\end{aligned}$$

$$\begin{aligned}J_L^\mu &= \bar{u} \gamma^\mu (1 - \gamma_5) d \\ J_R^\mu &= \bar{u} \gamma^\mu (1 + \gamma_5) d\end{aligned}$$

- Weight function: $w(z)\Pi(z)$
- $$w(s) = \frac{1}{s}, \frac{1}{s^2}$$

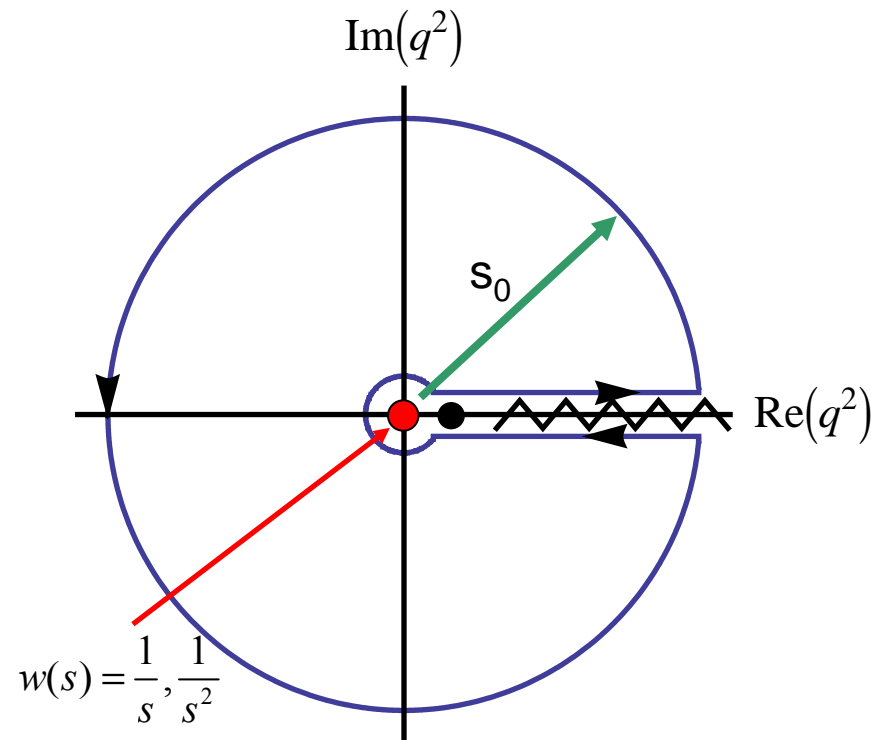
Analyticity of $w(z)\Pi(z)$ relates different regions of the C-plane.



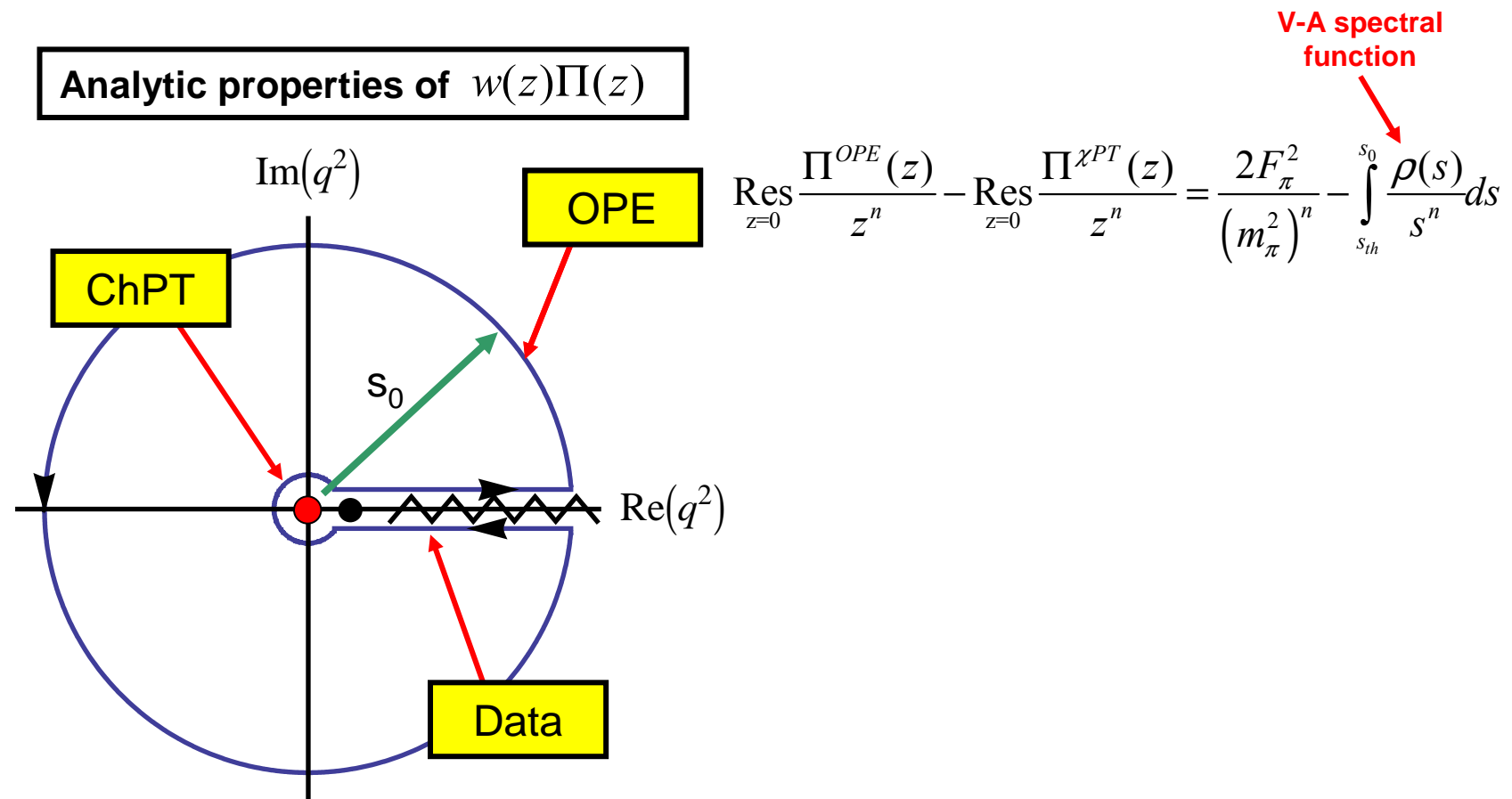
SUM RULE

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

Analytic properties of $w(z)\Pi(z)$

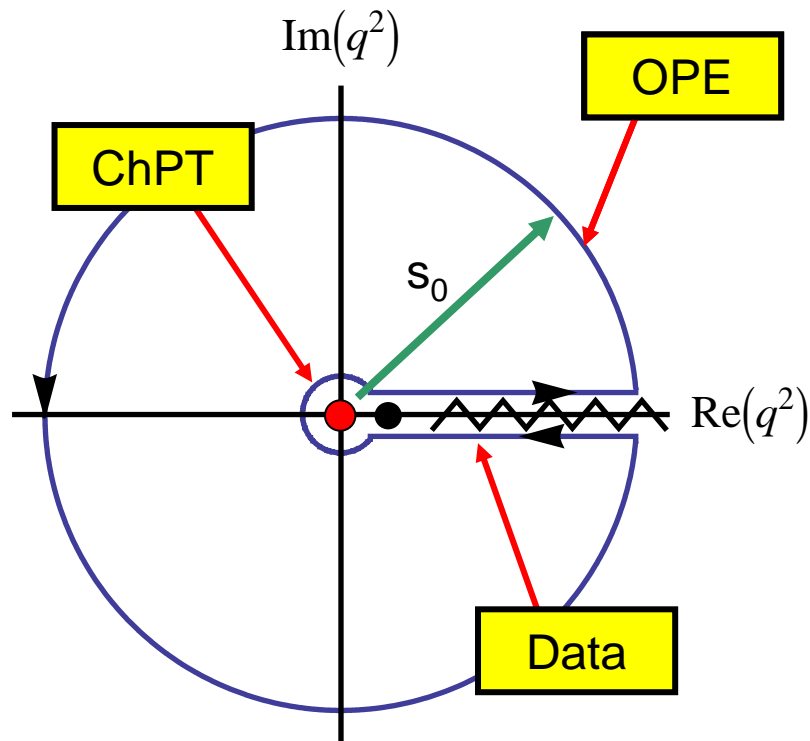


Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$



Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

Analytic properties of $w(z)\Pi(z)$

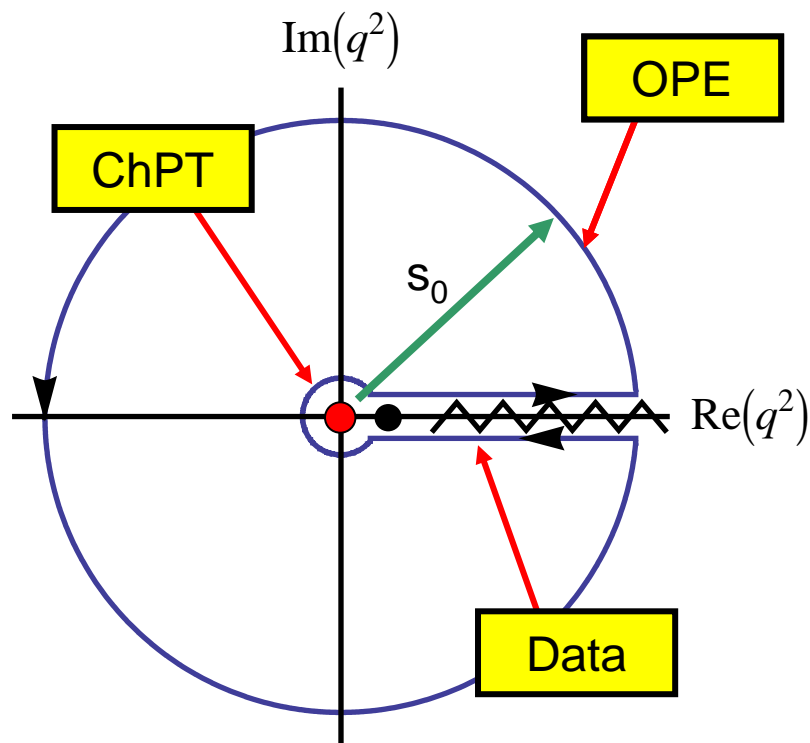


~~$$\text{Res}_{z=0} \frac{\Pi^{OPE}(z)}{z^n} - \text{Res}_{z=0} \frac{\Pi^{\chi PT}(z)}{z^n} = \frac{2F_\pi^2}{(m_\pi^2)^n} - \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$~~

$$\frac{d^{n-1}}{dz^{n-1}} \Pi^{\chi PT}(z) - \frac{2F_\pi^2}{(m_\pi^2)^n} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

Analytic properties of $w(z)\Pi(z)$



~~$$\text{Res}_{z=0} \frac{\Pi^{OPE}(z)}{z^n} - \text{Res}_{z=0} \frac{\Pi^{\chi PT}(z)}{z^n} = \frac{2F_\pi^2}{(m_\pi^2)^n} - \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$~~

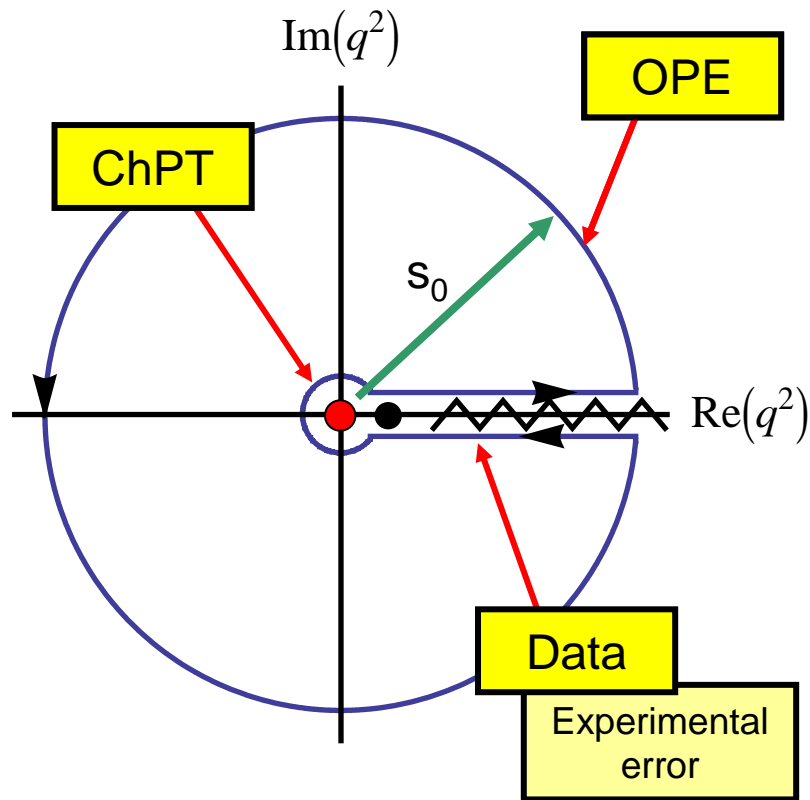
$$\frac{d^{n-1}}{dz^{n-1}} \Pi^{\chi PT}(z) - \frac{2F_\pi^2}{(m_\pi^2)^n} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$

$$\Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^2} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s} ds$$

$$\frac{d}{ds} \Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^4} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^2} ds$$

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

Analytic properties of $w(z)\Pi(z)$



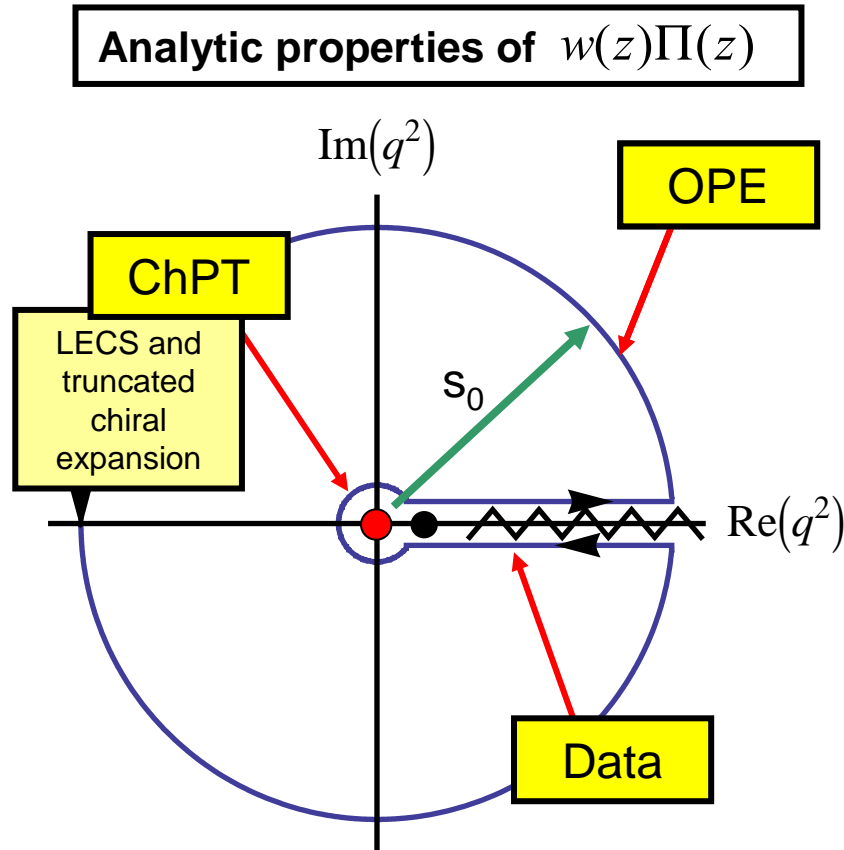
~~$$\text{Res}_{z=0} \frac{\Pi^{OPE}(z)}{z^n} - \text{Res}_{z=0} \frac{\Pi^{\chi PT}(z)}{z^n} = \frac{2F_\pi^2}{(m_\pi^2)^n} - \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$~~

$$\frac{d^{n-1}}{dz^{n-1}} \Pi^{\chi PT}(z) - \frac{2F_\pi^2}{(m_\pi^2)^n} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$

$$\Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^2} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s} ds$$

$$\frac{d}{ds} \Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^4} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^2} ds$$

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$



~~$$\text{Res}_{z=0} \frac{\Pi^{OPE}(z)}{z^n} - \text{Res}_{z=0} \frac{\Pi^{\chi PT}(z)}{z^n} = \frac{2F_\pi^2}{(m_\pi^2)^n} - \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$~~

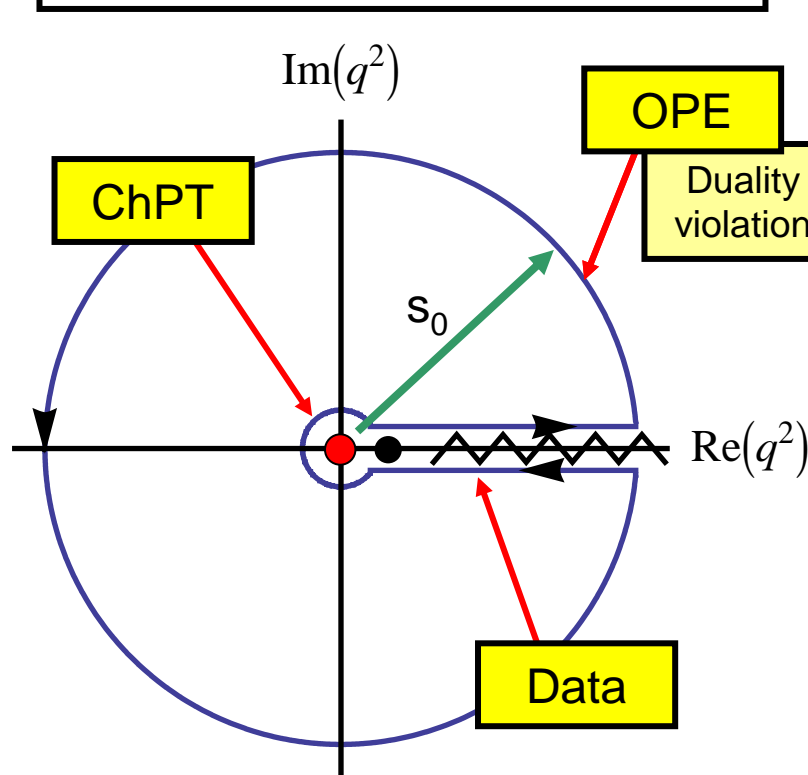
$$\frac{d^{n-1}}{dz^{n-1}} \Pi^{\chi PT}(z) - \frac{2F_\pi^2}{(m_\pi^2)^n} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$

$$\Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^2} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s} ds$$

$$\frac{d}{ds} \Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^4} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^2} ds$$

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

Analytic properties of $w(z)\Pi(z)$



$$\cancel{\text{Res}_{z=0} \frac{\Pi^{OPE}(z)}{z^n} - \text{Res}_{z=0} \frac{\Pi^{\chi PT}(z)}{z^n} = \frac{2F_\pi^2}{(m_\pi^2)^n} - \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds}$$

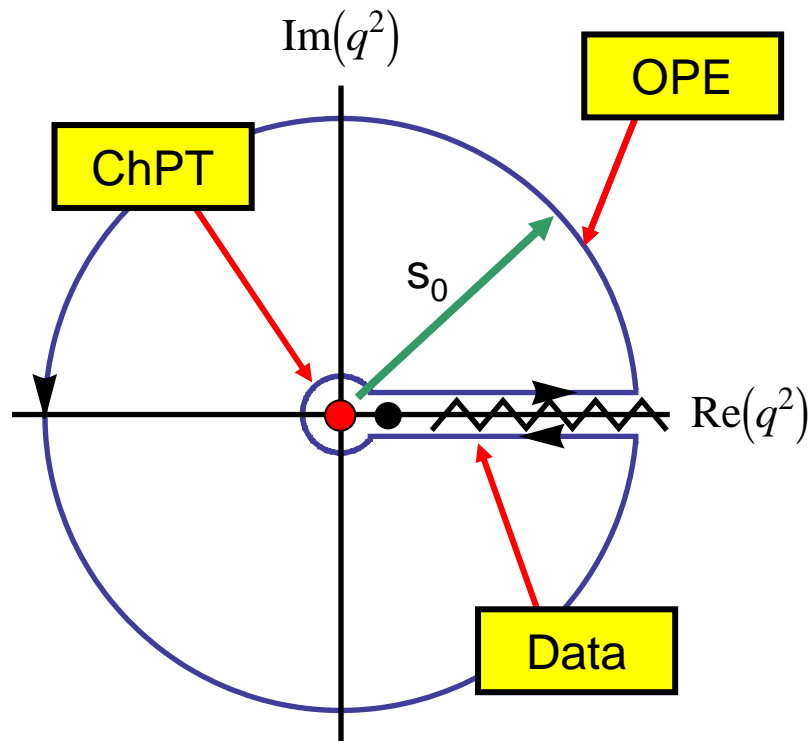
$$\frac{d^{n-1}}{dz^{n-1}} \Pi^{\chi PT}(z) - \frac{2F_\pi^2}{(m_\pi^2)^n} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$

$$\Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^2} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s} ds$$

$$\frac{d}{ds} \Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^4} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^2} ds$$

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

Analytic properties of $w(z)\Pi(z)$



~~$$\text{Res}_{z=0} \frac{\Pi^{OPE}(z)}{z^n} - \text{Res}_{z=0} \frac{\Pi^{\chi PT}(z)}{z^n} = \frac{2F_\pi^2}{(m_\pi^2)^n} - \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$~~

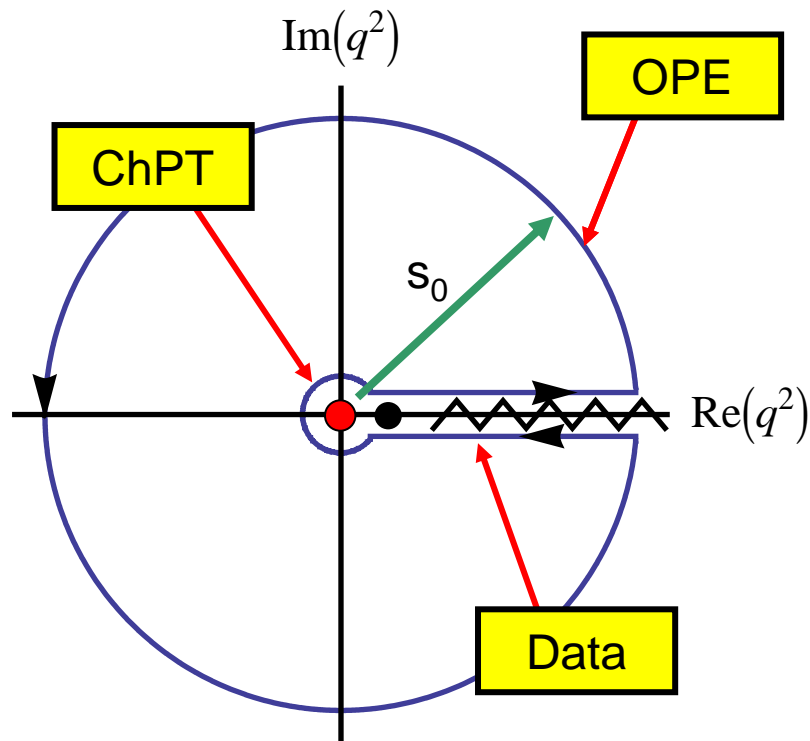
$$\frac{d^{n-1}}{dz^{n-1}} \Pi^{\chi PT}(z) - \frac{2F_\pi^2}{(m_\pi^2)^n} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$

$$\Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^2} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s} ds \equiv -8L_{10}^{eff}$$

$$\frac{d}{ds} \Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^4} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^2} ds \equiv 16C_{87}^{eff}$$

Approach: sum rule with $\Pi_{V-A}^{\mu\nu}(q)$

Analytic properties of $w(z)\Pi(z)$



~~$$\text{Res}_{z=0} \frac{\Pi^{OPE}(z)}{z^n} - \text{Res}_{z=0} \frac{\Pi^{\chi PT}(z)}{z^n} = \frac{2F_\pi^2}{(m_\pi^2)^n} - \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$~~

$$\frac{d^{n-1}}{dz^{n-1}} \Pi^{\chi PT}(z) - \frac{2F_\pi^2}{(m_\pi^2)^n} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^n} ds$$

$$\Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^2} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s} ds \quad \equiv -8L_{10}^{eff}$$

$$\frac{d}{ds} \Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^4} = \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^2} ds \quad \equiv 16C_{87}^{eff}$$

χ PT side

Data side

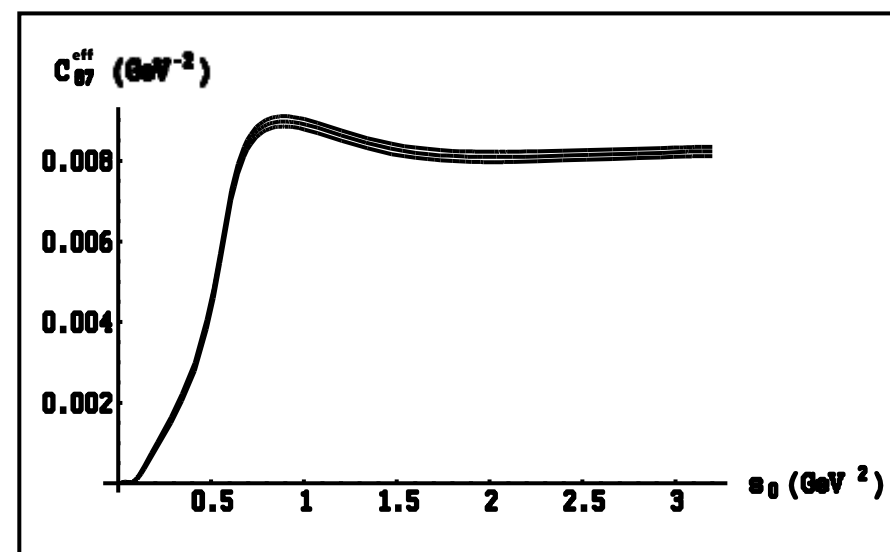
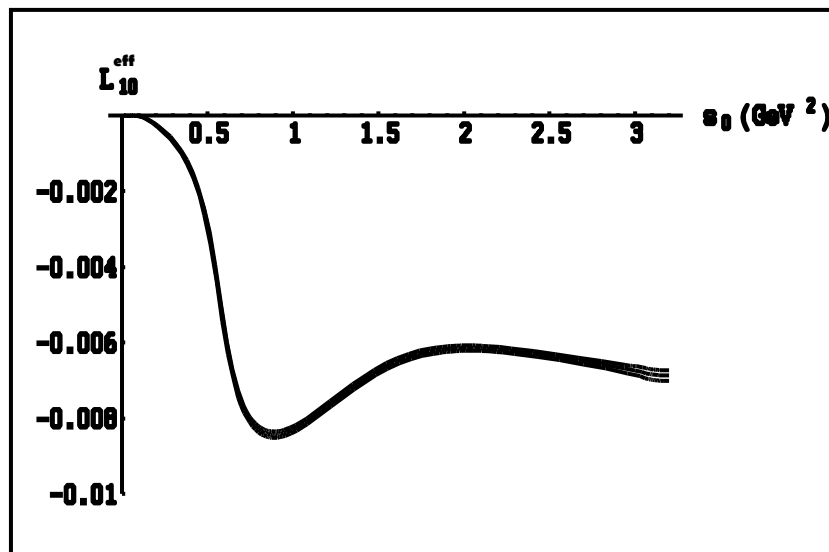
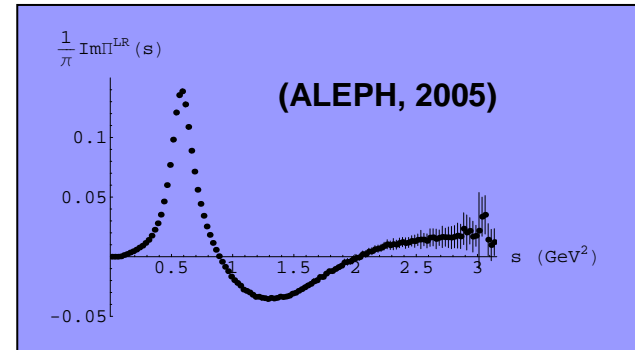
Data side

- Eff. parameters: determination.

$$-8L_{10}^{eff} \equiv \int_{s_{th}}^{s_0} \frac{\rho(s)}{s} ds$$

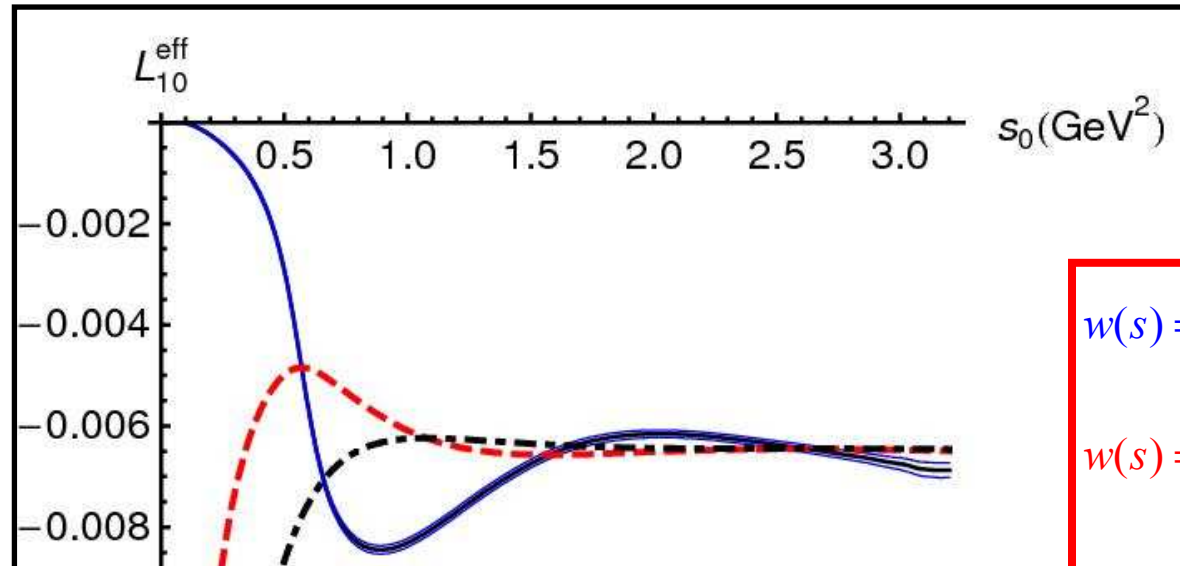
$$16C_{87}^{eff} \equiv \int_{s_{th}}^{s_0} \frac{\rho(s)}{s^2} ds$$

We can measure $\rho(s)$ from the hadronic tau decays...



Data side

- Eff. parameters: determination.



RESULTS:

$$L_{10}^{eff} = -(6.48 \pm 0.12) \cdot 10^{-3}$$

$$C_{87}^{eff} = +(8.18 \pm 0.14) \cdot 10^{-3} \text{ GeV}^{-2}$$

Good
accuracy!

$$w(s) = \frac{1}{s}$$

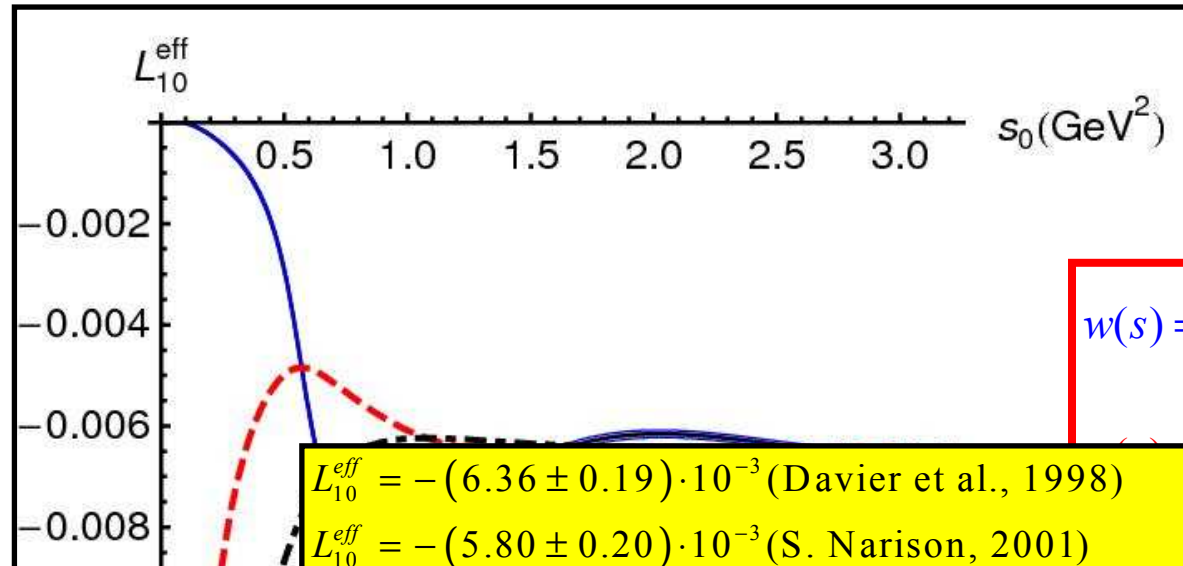
$$w(s) = \frac{1}{s} \left(1 - \frac{s}{s_0} \right)$$

$$w(s) = \frac{1}{s} \left(1 - \frac{s}{s_0} \right)^2$$

(F. Le Diberder & A. Pich, 1992,
K. Maltman, 1998,
Domínguez & Schilcher, 1999, ...)

Data side

- Eff. parameters: determination.



$$w(s) = \frac{1}{s}$$

$$L_{10}^{eff} = -(6.36 \pm 0.19) \cdot 10^{-3} \text{ (Davier et al., 1998)}$$

$$L_{10}^{eff} = -(5.80 \pm 0.20) \cdot 10^{-3} \text{ (S. Narison, 2001)}$$

$$L_{10}^{eff} = -(6.45 \pm 0.06) \cdot 10^{-3} \text{ (Domínguez et al., 2006)}$$

$$C_{87}^{eff} = +(8.10 \pm 0.12) \cdot 10^{-3} \text{ GeV}^{-2} \text{ (Domínguez \& Schilcher, 2004)}$$

RESULTS:

$$L_{10}^{eff} = -(6.48 \pm 0.12) \cdot 10^{-3}$$

$$C_{87}^{eff} = +(8.18 \pm 0.14) \cdot 10^{-3} \text{ GeV}^{-2}$$

Good accuracy!

(F. Le Diberder & A. Pich, 1992,
K. Maltman, 1998,
Domínguez & Schilcher, 1999, ...)

χ PT side

- Theoretical calculation of the effective parameters

$$L_{10}^{eff} = \frac{-1}{8} \left(\Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^2} \right) \equiv \frac{-1}{8} \bar{\Pi}^{\chi PT}(0) = f(L_{10}^r(\mu), \mu)$$

$$C_{87}^{eff} = \frac{1}{16} \left(\Pi^{\chi PT}{}'(0) + 2 \frac{F_\pi^2}{m_\pi^4} \right) \equiv \frac{1}{16} \bar{\Pi}^{\chi PT}{}'(0) = g(C_{87}^r(\mu), \mu)$$

- Using ChPT...

$$\bar{\Pi}^{\chi PT}(s) = -8L_{10}^r - 8B_V^{\pi\pi}(s) - 4B_V^{KK}(s) + O(p^6)$$

χ PT side

- Theoretical calculation of the effective parameters

$$L_{10}^{eff} = \frac{-1}{8} \left(\Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^2} \right) \equiv \frac{-1}{8} \bar{\Pi}^{\chi PT}(0) = f(L_{10}^r(\mu), \mu)$$

$$C_{87}^{eff} = \frac{1}{16} \left(\Pi^{\chi PT}'(0) + 2 \frac{F_\pi^2}{m_\pi^4} \right) \equiv \frac{1}{16} \bar{\Pi}^{\chi PT}'(0) = g(C_{87}^r(\mu), \mu)$$

- Using ChPT...

$$\begin{aligned} \bar{\Pi}^{\chi PT}(s) = & -8L_{10}^r - 8B_V^{\pi\pi}(s) - 4B_V^{KK}(s) \\ & + 16C_{87}^r s - 32m_\pi^2 (C_{61}^r - C_{12}^r - C_{80}^r) - 32(m_\pi^2 + 2m_K^2)(C_{62}^r - C_{13}^r - C_{81}^r) \\ & + \frac{16}{f_\pi^2} \left((2\mu_\pi + \mu_\pi)(L_9^r + 2L_{10}^r) - (2B_V^{\pi\pi}(s) + B_V^{KK}(s))sL_9^r \right) \\ & + G_{2L}(s, \mu) + O(p^8) \end{aligned}$$

(Amorós et al., 2000)
(Golowich & Kambor, 1995, 1998)

χ PT side

- Theoretical calculation of the effective parameters

$$L_{10}^{eff} = \frac{-1}{8} \left(\Pi^{\chi PT}(0) + 2 \frac{F_\pi^2}{m_\pi^2} \right) \equiv \frac{-1}{8} \bar{\Pi}^{\chi PT}(0) = f(L_{10}^r(\mu), \mu)$$

$$C_{87}^{eff} = \frac{1}{16} \left(\Pi^{\chi PT \prime}(0) + 2 \frac{F_\pi^2}{m_\pi^4} \right) \equiv \frac{1}{16} \bar{\Pi}^{\chi PT \prime}(0) = g(C_{87}^r(\mu), \mu)$$

- Therefore at $O(p^4)$...

$$L_{10}^{eff} = \frac{-1}{8} \bar{\Pi}^{\chi PT}(0) = L_{10}^r(\mu) + \left(\frac{\log \frac{m_K^2}{m_\pi^2}}{384\pi^2} + \frac{1 + \log \frac{m_\pi^2}{\mu^2}}{128\pi^2} \right)$$

$$L_{10}^r(m_\rho) = -(5.22 \pm 0.12) \cdot 10^{-3}$$

$$L_{10}^r(m_\rho) = -(5.13 \pm 0.19) \cdot 10^{-3}$$

(Davier et al., 1998)

χ PT side

- At $O(p^6)$...

$$L_{10}^{eff} = \frac{-1}{8} \bar{\Pi}^{\chi PT}(0)$$

$$C_{87}^{eff} = \frac{1}{16} \bar{\Pi}^{\chi PT}(0)$$

$$L_{10}^{eff} = L_{10}^r + \left(\frac{\log \frac{m_K^2}{m_\pi^2} + \frac{1 + \log \frac{m_\pi^2}{\mu^2}}{128\pi^2}}{384\pi^2} \right) +$$

$$+ 4m_\pi^2 (C_{61}^r - C_{12}^r - C_{80}^r) + 4(m_\pi^2 + 2m_K^2)(C_{62}^r - C_{13}^r - C_{81}^r)$$

$$- \frac{2}{f_\pi^2} (2\mu_\pi + \mu_\pi)(L_9 + 2L_{10}^r)$$

$$- \frac{1}{8} G_{2L}(0, \mu)$$

$$C_{87}^{eff} = \frac{1}{7680\pi^2} \left(\frac{1}{m_K^2} + \frac{2}{m_\pi^2} \right) +$$

$$+ C_{87}^r(\mu)$$

$$- \frac{2L_9(\mu)}{f_\pi^2} \left(\frac{\log \frac{m_K^2}{m_\pi^2} + \frac{1 + \log \frac{m_\pi^2}{\mu^2}}{128\pi^2}}{384\pi^2 f_\pi^2} + \frac{1}{128\pi^2} \right)$$

$$+ \frac{1}{16} G'_{2L}(0, \mu)$$

χ PT side

$$L_{10}^{eff} = \frac{-1}{8} \bar{\Pi}^{\chi PT}(0)$$

$$C_{87}^{eff} = \frac{1}{16} \bar{\Pi}^{\chi PT}(0)$$

- At $O(p^6)$...

$$L_{10}^{eff} = L_{10}^r + \left(\frac{\log \frac{m_K^2}{m_\pi^2} + \frac{1 + \log \frac{m_\pi^2}{\mu^2}}{128\pi^2}}{384\pi^2} \right) +$$

$$+ 4m_\pi^2 (C_{61}^r - C_{12}^r - C_{80}^r) + 4(m_\pi^2 + 2m_K^2)(C_{62}^r - C_{13}^r - C_{81}^r)$$

$$- \frac{2}{f_\pi^2} (2\mu_\pi + \mu_\pi)(L_9 + 2L_{10}^r)$$

$$- \frac{1}{8} G_{2L}(0, \mu)$$

Experimental error LEC's error

$$L_{10}^r(m_\rho) = -(4.06 \pm 0.08 \pm 0.39) \cdot 10^{-3}$$

$$= -(4.06 \pm 0.40) \cdot 10^{-3}$$

$$C_{87}^{eff} = \frac{1}{7680\pi^2} \left(\frac{1}{m_K^2} + \frac{2}{m_\pi^2} \right) +$$

$$+ C_{87}^r(\mu)$$

$$- \frac{2L_9(\mu)}{f_\pi^2} \left(\frac{\log \frac{m_K^2}{m_\pi^2} + \frac{1 + \log \frac{m_\pi^2}{\mu^2}}{128\pi^2}}{384\pi^2 f_\pi^2} + \frac{1}{128\pi^2} \right)$$

$$+ \frac{1}{16} G'_{2L}(0, \mu)$$

$$C_{87}^r(m_\rho) = +(4.89 \pm 0.14 \pm 0.13) \cdot 10^{-3} \text{ GeV}^2$$

$$= +(4.89 \pm 0.19) \cdot 10^{-3} \text{ GeV}^2$$

Comparisons

- L_{10} : extractions from the data:

$$L_{10}^r(m_\rho) = -(4.06 \pm 0.40) \cdot 10^{-3}$$

$$L_{10}^r(m_\rho) = -(5.13 \pm 0.19) \cdot 10^{-3} \text{ (From inclusive } \tau \text{ data, Davier et al., 1998)}$$

$$L_{10}^r(m_\rho) = -(4.5 \pm 0.4) \cdot 10^{-3} \text{ (From } \pi \rightarrow e\nu\gamma \text{ \& } \langle r^2 \rangle_\pi, 2008)$$

- L_{10} : theoretical estimations:

μ -sensitive
(NLO calculation)

Not μ -sensitive
(LO calculation)

$$L_{10}^r(m_\rho) = -(4.4 \pm 0.9) \cdot 10^{-3} \text{ (Pich et al., 2008)}$$

$$L_{10}^r(m_\rho) = -(5.7 \pm 1.4) \cdot 10^{-3} \text{ (Ecker et al., 1989)}$$

(updated)

25%
assigned
error

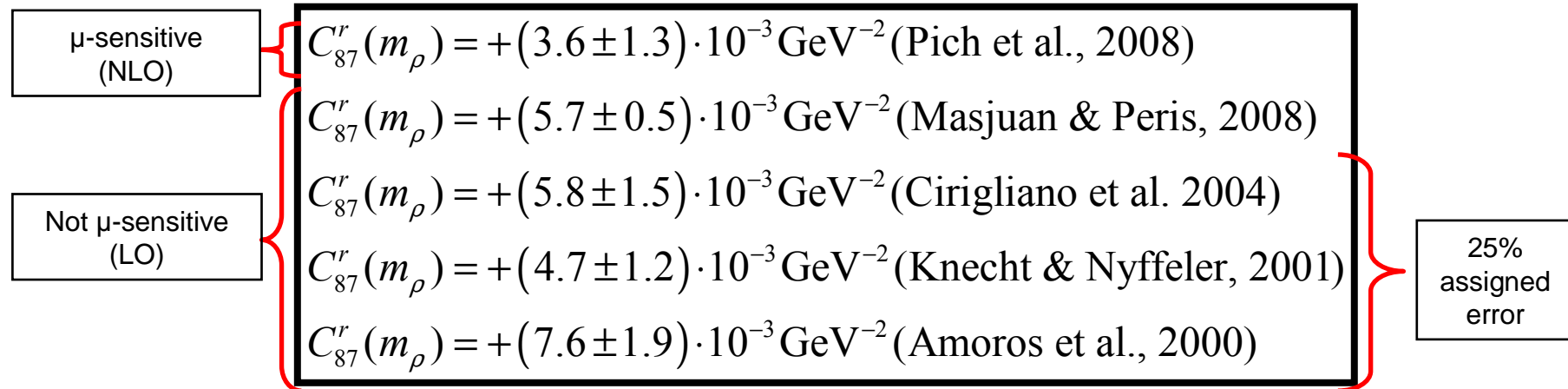
Lattice...

$$L_{10}^r(m_\rho) = -(5.2 \pm 0.5) \cdot 10^{-3} \text{ (Shintani et al., 2008)}$$

Comparisons

- C_{87} : With theoretical (model dependent) estimations:

$$C_{87}^r(m_\rho) = +(4.89 \pm 0.19) \cdot 10^{-3} \text{ GeV}^{-2}$$



L_9^r , \bar{l}_5 and \bar{l}_6

- SU(3) χ PT LECs $\xleftrightarrow{O(p^6)}$ SU(2) χ PT LECs (J. Gasser et al., 2007)

$$L_{10}^r(m_\rho) \Rightarrow \bar{l}_5 = +12.24 \pm 0.29 \leftarrow O(p^6)$$

- L_9 & l_6 determination:

$$L_9^r(m_\rho) = +(5.93 \pm 0.43) \cdot 10^{-3} \quad \text{OK!}$$

(From $\langle r^2 \rangle_\pi$, Bijens & Talavera, 2002)

$$\begin{aligned} (L_9^r + L_{10}^r)(m_\rho) = -(1.44 \pm 0.08) \cdot 10^{-3} &\Rightarrow L_9^r(m_\rho) = +(5.50 \pm 0.41) \cdot 10^{-3} \leftarrow O(p^6) \\ \bar{l}_6 - \bar{l}_5 = +(2.98 \pm 0.33) \cdot 10^{-3} &\Rightarrow \bar{l}_6 = +15.22 \pm 0.44 \leftarrow O(p^6) \end{aligned}$$

(From $\pi \rightarrow e\nu\gamma$,

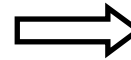
J. Bijens & P. Talavera, 1997,

R. Unterdorfer & H. Pichl, 2008)

Conclusions

- We have calculated L_{10} with recent inclusive tau data including $O(p^6)$ terms in the chiral expansion, obtaining

$$L_{10}^r(m_\rho) = -(4.06 \pm 0.40) \cdot 10^{-3}$$



$$\begin{aligned} L_9^r(m_\rho) &= +(5.50 \pm 0.41) \cdot 10^{-3} \\ \bar{l}_5 &= +12.24 \pm 0.29 \\ \bar{l}_6 &= +15.22 \pm 0.44 \end{aligned}$$

- The same approach has allowed us to calculate the first determination from data of C_{87} (up to $O(p^6)$ terms) obtaining

$$C_{87}^r(m_\rho) = +(4.89 \pm 0.19) \cdot 10^{-3} \text{ GeV}^{-2}$$

... in good agreement with different theoretical predictions.

Backup

Operator Product Expansion

(Wilson, 1967):

$$\Pi_{V-A}(q^2) = \Pi_{V-A}^{pert}(q^2) + \frac{\mathcal{O}_2}{-q^2} + \frac{\mathcal{O}_4}{(-q^2)^2} + \frac{\mathcal{O}_6}{(-q^2)^3} + \dots$$

The ChPT Lagrangian

$$\begin{aligned} \mathcal{L}^{\text{ChPT}} = & \frac{F^2}{4} \langle u_\mu u^\mu + \chi_+ \rangle + \\ & + \dots + \frac{1}{4} L_{10} \langle f_{+\mu\nu} f_+^{\mu\nu} - f_{-\mu\nu} f_-^{\mu\nu} \rangle + \dots + \\ & + \dots + C_{87} \langle \nabla_\rho f_{-\mu\nu} \nabla^\rho f_-^{\mu\nu} \rangle + \dots + O(p^8) \end{aligned}$$

$$\begin{aligned} u_\mu &= i \left[u^\dagger (\partial_\mu - i r_\mu) u - u (\partial_\mu - i l_\mu) u^\dagger \right] \\ \chi_\pm &= u^\dagger \chi u^\dagger \pm u \chi^\dagger u \\ f_\pm^{\mu\nu} &= u F_L^{\mu\nu} u^\dagger \pm u^\dagger F_R^{\mu\nu} u \end{aligned}$$

Meson fields...

$$u(\phi) = e^{\frac{i}{\sqrt{2}F}\phi}$$

$$\phi = \frac{1}{\sqrt{2}} \sum_{i=1}^8 \lambda_i \phi_i = \begin{pmatrix} \frac{1}{\sqrt{2}} \pi^0 + \frac{1}{\sqrt{6}} \eta_8 & \pi^+ & K^+ \\ \pi^- & \frac{-1}{\sqrt{2}} \pi^0 + \frac{1}{\sqrt{6}} \eta_8 & K^0 \\ K^- & \bar{K}^0 & \frac{-2}{\sqrt{6}} \eta_8 \end{pmatrix}$$

External fields...

$$\chi = 2B_0 (s + ip)$$

$$F_R^{\mu\nu} = \partial^\mu r^\nu - \partial^\nu r^\mu - i[r^\mu, r^\nu]$$

$$F_L^{\mu\nu} = \partial^\mu l^\nu - \partial^\nu l^\mu - i[l^\mu, l^\nu]$$

The explicit breaking of chiral symmetry through quark masses can be added by...

$$s = \mathfrak{M} = \begin{pmatrix} m_u & 0 & 0 \\ 0 & m_d & 0 \\ 0 & 0 & m_s \end{pmatrix}$$